

3) *Doubly Exponential Probability Density Function*: The calculation of  $\phi^2(\alpha, k_e)$  where

$$k_e(x) = \frac{1}{2} \exp(-|x|)$$

is similar to the calculation of  $\phi^2(\alpha, k_n)$ .

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## Bounds on the Sum of Element Sensitivity Magnitudes for Network Functions

M. N. S. SWAMY, CHAMPA BHUSHAN, AND K. THULASIRAMAN

**Abstract**—Bounds on the sum of element sensitivity magnitudes for transfer immittances, transfer voltage, and current ratios are established for networks consisting of resistors, capacitors, inductors, and gyrators. For RC and LC networks these bounds are given in terms of relevant driving-point functions and their frequency sensitivities. Further, for two-element-kind networks, bounds are given for the quadratic sensitivity index  $\phi$ .

### I. INTRODUCTION

Recently, Smith [1] has obtained, using Vratsano's theorem, bounds on the sum of element sensitivity magnitudes for driving-point impedances (DPI) and for networks consisting of resistors, inductors, capacitors, and gyrators. In this paper we use the adjoint network approach [2] to establish for the same class of networks bounds on the sum of element sensitivity magnitudes for transfer immittances as well as transfer voltage and current ratios.

### II. BOUNDS ON THE SUM OF ELEMENT SENSITIVITY MAGNITUDES FOR NETWORK FUNCTIONS

#### A. Transfer Impedance $Z_{12}$

Consider a network  $N$  consisting of resistors, inductors, capacitors, and gyrators. Let  $N'$  denote the adjoint of  $N$ . Let  $Z_{12}$  be the transfer impedance between ports 1 and 2. Consider an element  $z_n$ . It is known from [1] that

$$\frac{\partial Z_{12}}{\partial z_n} = \frac{I_n I_n'}{I_0^2} \quad (1)$$

where  $I_n (I_n')$  is the current through the element  $z_n$  when port 2(1) of  $N(N')$  is excited with a current source of value  $I_0$  and port 1(2) is open circuited.  $|I_0|$  refers to the effective value of  $I_0$ . It follows from (1) that

$$|S_{z_n}^{Z_{12}}| = \frac{|I_n| |I_n'|}{|I_0|^2} \times \frac{|z_n|}{|Z_{12}|} \quad (2)$$

where  $S_{z_n}^{Z_{12}}$  is the sensitivity of  $Z_{12}$  with respect to  $z_n$ . It may be mentioned that the sensitivity of any function  $F$  with respect to a parameter  $p$  is defined as

$$S_p^F = \frac{p}{F} \frac{\partial F}{\partial p} \quad (3)$$

If  $z_n$  corresponds to a resistor  $r$ , then from (1) we get that

$$|S_r^{Z_{12}}| = \frac{|I_r| |I_r'|}{|I_0|^2} \times \frac{r}{|Z_{12}|} = \frac{\sqrt{P_r P_r'}}{|I_0|^2 |Z_{12}|} \quad (4)$$

where  $P_r (P_r')$  denotes the average power dissipated in the resistor  $r$  when port 2(1) of  $N(N')$  is excited with a current source of value  $I_0$  and port 1(2) is open circuited.

If  $z_n$  represents an inductor  $l$  or a capacitor  $c$ , then from (2) we get the following:

$$|S_l^{Z_{12}}| = \frac{2\omega \sqrt{W_l W_l'}}{|I_0|^2 |Z_{12}|} \quad (5a)$$

and

$$|S_c^{Z_{12}}| = \frac{2\omega \sqrt{W_c W_c'}}{|I_0|^2 |Z_{12}|} \quad (5b)$$

where  $W_l (W_c)$  and  $W_l' (W_c')$  denote the average energy stored in the inductor (capacitor) in the network  $N$  and  $N'$ , respectively. It should be noted that  $N$  and  $N'$  are terminated at the ports as described earlier.

Let  $P_R (P_R')$  denote the total average power dissipated in  $N(N')$ . Then

$$P_R = \sum_r P_r \quad P_R' = \sum_r P_r' \quad (6)$$

Let at frequency  $\omega$

$$Z_{22} = R_{22} + jX_{22} \quad Z_{11}' = R_{11}' + jX_{11}' \quad (7)$$

Using (4), (6), and (7), along with Schwarz's inequality, we get the following:

$$\sum_r |S_r^{Z_{12}}| = \frac{\sum_r \sqrt{P_r P_r'}}{|I_0|^2 |Z_{12}|} \leq \frac{\sqrt{P_R P_R'}}{|I_0|^2 |Z_{12}|} = \frac{\sqrt{R_{22} R_{11}'}}{|Z_{12}|} = \frac{\sqrt{R_{22} R_{11}}}{|Z_{12}|} \quad (8)$$

The last step in (8) follows from the relation  $Z_{11}' = Z_{11}$  [3].

It may be noted that in the case of a DPI, (8) will reduce to [1, eq. 45]. Further, if  $N$  consists of  $r$ ,  $l$ , and  $c$  elements only, then for a DPI, (8) will reduce to [1, eq. 11] since the adjoint of  $N$  is itself if  $N$  is reciprocal.

Let

$$W_C = \sum_c W_c \quad W_C' = \sum_c W_c' \quad W_L = \sum_l W_l$$

$$W_L' = \sum_l W_l' \quad W = W_C + W_L \quad W' = W_C' + W_L' \quad (9)$$

Using (5), (7), and (9), along with Schwarz's inequality we get the following:

$$\sum_l |S_l^{Z_{12}}| + \sum_c |S_c^{Z_{12}}|$$

$$= \frac{2\omega \left( \sum_l \sqrt{W_l W_l'} + \sum_c \sqrt{W_c W_c'} \right)}{|I_0|^2 |Z_{12}|} \leq \frac{2\omega \sqrt{W W'}}{|I_0|^2 |Z_{12}|} \quad (10)$$

If the network  $N$  is reciprocal, then using (5), [1, eq. 12] for a DPI can be easily established in a similar manner.

It is known [4] that for the network under consideration

$$\sum_l S_l^F + \sum_c S_c^F = S_w^F \quad (11)$$

where  $F$  is any network function. Hence we get

$$\sum_l |S_l^F| + \sum_c |S_c^F| \geq |S_w^F| \quad (12)$$

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From (10) and (12) we get the following:

$$|S_{\omega}^F| \leq \sum_i |S_i^{Z_{12}}| + \sum_c |S_c^{Z_{12}}| \leq \frac{2\omega\sqrt{WW'}}{|I_0|^2 |Z_{12}|}. \quad (13)$$

If the network consists of  $r$  and  $c$  elements only, then it is known that [4], [5]

$$\sum_r S_r^{Z_{12}} = 1 + S_{\omega}^{Z_{12}} \quad \sum_c S_c^{Z_{12}} = S_{\omega}^{Z_{12}}. \quad (14)$$

Hence, on taking magnitudes we have

$$\sum_r |S_r^{Z_{12}}| + \sum_c |S_c^{Z_{12}}| \geq |1 + 2S_{\omega}^{Z_{12}}|. \quad (15)$$

Further, for an  $RC$  network we can show starting from (5b) that

$$\sum_c |S_c^{Z_{12}}| \leq (\sqrt{|X_{11}X_{22}}|)/|Z_{12}|. \quad (16)$$

Then from (8), (15), and (16) we obtain the following inequality:

$$\begin{aligned} |1 + 2S_{\omega}^{Z_{12}}| &\leq \sum_r |S_r^{Z_{12}}| + \sum_c |S_c^{Z_{12}}| \\ &\leq \frac{\sqrt{R_{11}R_{22}} + \sqrt{|X_{11}X_{22}}|}{|Z_{12}|}. \end{aligned} \quad (17)$$

An inequality similar to (17) can be obtained easily for an  $RL$  network.

Let the network  $N$  consist of only inductors and capacitors. It is known that [6]

$$W = W_C + W_L = \frac{1}{2} |I_0|^2 \frac{\partial X_{22}}{\partial \omega}$$

and

$$W' = W_L' + W_C' = \frac{1}{2} |I_0|^2 \frac{\partial X_{11}}{\partial \omega}. \quad (18)$$

Hence from (10) and (18) we get the following inequality:

$$\begin{aligned} \sum_i |S_i^{Z_{12}}| + \sum_c |S_c^{Z_{12}}| &\leq \frac{\omega}{|Z_{12}|} \sqrt{\frac{\partial X_{22}}{\partial \omega} \frac{\partial X_{11}}{\partial \omega}} \\ &= \sqrt{(|S_{\omega}^{Z_{22}}| |S_{\omega}^{Z_{11}}| |Z_{11}| |Z_{22}|)/|Z_{12}|}. \end{aligned} \quad (19)$$

Thus from (13) and (19) for an  $LC$  network we get the following inequality:

$$\begin{aligned} |S_{\omega}^{Z_{12}}| &\leq \sum_i |S_i^{Z_{12}}| + \sum_c |S_c^{Z_{12}}| \\ &\leq \sqrt{(|S_{\omega}^{Z_{22}}| |S_{\omega}^{Z_{11}}| \times |Z_{11}| |Z_{22}|)/|Z_{12}|}. \end{aligned} \quad (20)$$

This completes our discussion of the bounds on the sum of element sensitivity magnitudes for transfer impedance functions.

#### B. Transfer Admittance $Y_{12}$

Let  $Y_{12}$  denote the transfer admittance between ports 1 and 2 of a network  $N$ . Let port 2(1) of  $N(N')$  be excited with a voltage source of value  $V_0$  and port 1(2) be short circuited. Let

$$Y_{22} = G_{22} + jB_{22} \quad Y_{11}' = G_{11}' + jB_{11}'$$

where  $Y_{22}(Y_{11}')$  is the short-circuit driving-point admittance across port 2(1) of  $N(N')$ . The quantities  $W_i$ ,  $W_c$ ,  $W_i'$ ,  $W_c'$ , etc. are the same as in the case of transfer impedances, except that  $N$  and  $N'$  are now terminated as mentioned above.

With these definitions all inequalities and equations from (4) to (20) hold true if  $Z_{12}$ ,  $I_0$ ,  $R_{11}$ ,  $R_{22}$ ,  $X_{11}$ , and  $X_{22}$  are replaced, respectively, by  $Y_{12}$ ,  $V_0$ ,  $G_{11}$ ,  $G_{22}$ ,  $B_{11}$ , and  $B_{22}$ , and in (14), (15), and (17) the quantity  $(2S_{\omega}^{Z_{12}} + 1)$  is replaced by  $(2S_{\omega}^{Y_{12}} - 1)$ .

#### C. Transfer Voltage Ratio $T_V$

Let  $T_V$  be the transfer voltage ratio between ports 2 and 1 of a network  $N$ . Let port 1(2) of  $N(N')$  be excited with a voltage (current) source of value  $V_1(I_2)$  and port 2(1) of  $N(N')$  be open circuited (short circuited). Let  $W(W')$  denote the average energy stored in the reac-

TABLE I  
BOUNDS ON THE SUM OF ELEMENT SENSITIVITY  
MAGNITUDES FOR  $T_V$

Network Type	Bounds
$R, L, C,$ and Gyrators	$\sum_r  S_r^{T_V}  \leq (\sqrt{G_{11}R_{22}}/ T_V )$ $ S_{\omega}^{T_V}  \leq \sum_i  S_i^{T_V}  + \sum_c  S_c^{T_V} $ $\leq (2\omega\sqrt{WW'})/ V_1  I_2  T_V $
$R$ and $C$	$ S_{\omega}^{T_V}  \leq \sum_r  S_r^{T_V}  \leq \sqrt{G_{11}R_{22}}/ T_V $ $ S_{\omega}^{T_V}  \leq \sum_c  S_c^{T_V}  \leq \sqrt{B_{11}X_{22}}/ T_V $
$L$ and $C$	$ S_{\omega}^{T_V}  \leq \sum_i  S_i^{T_V}  + \sum_c  S_c^{T_V} $ $\leq (1/ T_V )\sqrt{ S_{\omega}^{Y_{11}}  S_{\omega}^{Z_{22}}  Y_{11}  Z_{22}}}$

tive elements of  $N(N')$  when it is terminated as above. Let for  $N$  and  $N'$

$$Y_{11} = 1/Z_{11} = G_{11} + jB_{11} \quad \text{and} \quad Z_{22}' = 1/Y_{22}' = R_{22}' + jX_{22}'.$$

Following the same approach as used in the case of  $Z_{12}$ , we can establish bounds on the sum of element sensitivity magnitudes for  $T_V$ . These results are summarized in Table I.

#### D. Transfer Current Ratio $T_c$

Let  $T_c$  denote the transfer current ratio between ports 2 and 1 of a network  $N$ . Let  $W(W')$  denote the total average energy stored in the reactive elements of  $N(N')$  when port 1(2) of  $N(N')$  is excited with a current (voltage) source of value  $I_1(V_2)$  and port 2(1) of  $N(N')$  is short circuited (open circuited). Let for  $N$  and  $N'$

$$Z_{11} = 1/Y_{11} = R_{11} + jX_{11} \quad \text{and} \quad Y_{22}' = 1/Z_{22}' = G_{22}' + jB_{22}'.$$

Observing that the transfer current ratio between ports 2 and 1 of  $N$  is the same as the transfer voltage ratio between ports 1 and 2 of  $N'$ , it can be shown that all the inequalities of Table I will hold true if  $I_2$ ,  $V_1$ ,  $G_{11}$ ,  $B_{11}$ ,  $R_{22}$ , and  $X_{22}$  are replaced, respectively, by  $I_1$ ,  $V_2$ ,  $G_{22}$ ,  $B_{22}$ ,  $R_{11}$ , and  $X_{11}$ .

### III. BOUNDS ON QUADRATIC SENSITIVITY INDEX

A commonly used criterion of performance is the quadratic sensitivity index  $\phi$  given by [7]

$$\phi = \sum_{x_i} |S_{x_i}^F|^2.$$

In the previous discussion we have obtained for two-element-kind networks bounds for the sum of magnitude sensitivities with respect to all the elements of the network. Using these results we can obtain bounds for the function  $\phi$ . For an  $RC$  network, we get from (8) and (16) the following:

$$\sum_r |S_r^{Z_{12}}|^2 \leq R_{11}R_{22}/|Z_{12}|^2 \quad (21a)$$

$$\sum_c |S_c^{Z_{12}}|^2 \leq X_{11}X_{22}/|Z_{12}|^2. \quad (21b)$$

Hence

$$\phi = \sum_r |S_r^{Z_{12}}|^2 + \sum_c |S_c^{Z_{12}}|^2 \leq \frac{R_{11}R_{22}}{|Z_{12}|^2} + \frac{X_{11}X_{22}}{|Z_{12}|^2}. \quad (22)$$

Using (14) and the Chebyshev inequality

$$\left(\sum_{i=1}^n a_i\right)^2/n \leq \sum_{i=1}^n a_i^2 \quad (23)$$

we can show that

$$\phi \geq \frac{1}{n_1} |1 + S_{\omega}^{Z_{12}}|^2 + \frac{1}{n_2} |S_{\omega}^{Z_{12}}|^2 \quad (24)$$

TABLE II  
BOUNDS ON QUADRATIC SENSITIVITY INDEX  $\phi$

$F$	RC Network <sup>a</sup>	LC Network <sup>b</sup>
$Z_{ij}$	$\frac{1}{n_1}  1 + S_{\omega} F ^2 + \frac{1}{n_2}  S_{\omega} F ^2 \leq \phi \leq (R_{11}R_{22} +  X_{11}  X_{22} )/ F ^2$	$\frac{1}{4n_1}  S_{\omega} F - 1 ^2 + \frac{1}{4n_2}  S_{\omega} F + 1 ^2 \leq \phi \leq ( S_{\omega} Z_{22}  S_{\omega} Z_{11}  Z_{11}  Z_{22} )/ F ^2$
$Y_{ij}$	$\frac{1}{n_1}  -1 + S_{\omega} F ^2 + \frac{1}{n_2}  S_{\omega} F ^2 \leq \phi \leq (G_{11}G_{22} +  B_{11}  B_{22} )/ F ^2$	$\frac{1}{4n_1}  S_{\omega} F + 1 ^2 + \frac{1}{4n_2}  S_{\omega} F - 1 ^2 \leq \phi \leq ( S_{\omega} Y_{22}  S_{\omega} Y_{11}  Y_{11}  Y_{22} )/ F ^2$
$T_V$	$\left(\frac{1}{n_1} + \frac{1}{n_2}\right)  S_{\omega} F ^2 \leq \phi \leq (G_{11}R_{22} +  B_{11}  X_{22} )/ F ^2$	$\frac{1}{4} \left(\frac{1}{n_1} + \frac{1}{n_2}\right)  S_{\omega} F ^2 \leq \phi \leq ( S_{\omega} Y_{11}  S_{\omega} Z_{22}  Y_{11}  Z_{22} )/ F ^2$
$T_c$	$\left(\frac{1}{n_1} + \frac{1}{n_2}\right)  S_{\omega} F ^2 \leq \phi \leq (G_{22}R_{11} +  B_{22}  X_{11} )/ F ^2$	$\frac{1}{4} \left(\frac{1}{n_1} + \frac{1}{n_2}\right)  S_{\omega} F ^2 \leq \phi \leq ( S_{\omega} Y_{22}  S_{\omega} Z_{11}  Y_{22}  Z_{11} )/ F ^2$

<sup>a</sup>  $n_1$  = number of resistors;  $n_2$  = number of capacitors.  
<sup>b</sup>  $n_1$  = number of capacitors;  $n_2$  = number of inductors.

where  $n_1$  is the number of resistors and  $n_2$  is the number of capacitors in the network.

Similar bounds can be obtained for other network functions of two-element-kind networks. These are summarized in Table II.

#### IV. CONCLUSIONS

For networks containing  $R$ ,  $L$ , and  $C$  elements and gyrators, bounds have been obtained for the sum of magnitude of sensitivities of network functions with respect to  $R$ ,  $C$ , and  $L$ . In the case of two-element-kind networks these bounds have been expressed in terms of frequency sensitivities and the real and imaginary parts of relevant driving-point functions. Further, bounds have been obtained for the quadratic sensitivity index  $\phi$  for the case of two-element-kind networks.

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### Simple Bounds on the Variance of System Functions of Two-Element-Kind Networks Having Randomly Chosen Elements

J. C. GIGUERE

**Abstract**—The elements of the networks under consideration are assumed to have small statistically independent random variations about their nominal values. With these assumptions, it is shown that the variance of a given system function can be related to its sensitivity with respect to the network elements. Recent results in sensitivity analysis are then used to establish simple upper bounds on the variance of two-port network functions.

#### INTRODUCTION

Let  $F$  be a system function of an arbitrary network composed of resistors, capacitors, and inductors having nominal values  $r_i$ ,  $c_i$ , and  $l_i$ . Let  $\Delta x_i$  be a small random fluctuation of a given element about its nominal value  $x_i$  such that  $E[\Delta x_i] = 0$ . Then the fluctuation  $\Delta F_{x_i}$  of the system function about its nominal value may be approximated by

$$\Delta F_{x_i} \approx \frac{\partial F}{\partial x_i} \Delta x_i.$$

Now defining the normalized fluctuation of the element as  $\epsilon_i = \Delta x_i/x_i$ , the normalized fluctuation  $V_{x_i} F \triangleq \Delta F_{x_i}/F$  of the system function  $F$  may be expressed as

$$V_{x_i} F = \frac{\partial F}{\partial x_i} \frac{x_i}{F} \epsilon_i = S_{x_i} F \epsilon_i.$$

For the whole set of elements,<sup>1</sup> we have

$$V_x F = \sum_i S_{x_i} F \epsilon_i, \quad x = \{x_1, x_2, \dots, x_n\}. \quad (3)$$

It is evident that  $E[V_x F] = 0$ . Assuming that the random variables  $\epsilon_i$  are statistically independent, we have that  $E[\epsilon_i \epsilon_j] = 0$ ,  $i \neq j$ , and hence [1]

$$E|V_x F|^2 = \sum_i |S_{x_i}(j\omega)|^2 \sigma_{\epsilon_i}^2, \quad \sigma_{\epsilon_i}^2 = E[\epsilon_i^2]. \quad (4)$$

We now note that all those elements which have been chosen according to the same probability law (e.g., uniformly distributed 10-percent resistors) will have identical normalized variances. In such cases we will write  $\sigma_{\epsilon_i}^2 = \sigma_x^2$  where now  $x$  denotes the class of elements defined by

$$x = \{x_i: E[\Delta x_i/x_i]^2 = \sigma_x^2\}.$$

Equation (4) now becomes

$$E|V_x F|^2 = \sigma_x^2 \sum_i |S_{x_i} F|^2. \quad (5)$$

In what follows we will assume that the sets of resistors, capacitors, and inductors each constitute such a class.<sup>2</sup>

#### Bounds on the Variance of $V_x F$

We will now make use of recent results in sensitivity analysis [2], [3], to formulate upper bounds on  $E|V_x F|^2$ .

**Driving Point Functions:** Smith [2] has shown for driving point functions composed of resistors, capacitors, inductors, and ideal transformers that

$$\sum_i |S_{r_i} Z_0| = \frac{R_0}{|Z_0|}, \quad Z_0 = R_0 + jX_0 \quad (6a)$$

where  $S_{r_i} Z_0$  is the sensitivity of the impedance  $Z_0$  with respect to the

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<sup>1</sup> The inequalities obtained will still be valid for the more general case if we let  $\sigma_x^2 = \max_i \{\sigma_{\epsilon_i}^2\}$ .  
<sup>2</sup> A similar measure has been formulated by Rosenblum and Ghauis [4].